

# Soliton Transmission Capacity of Vertical Cavity Surface Emitting Lasers (VCSELs) Degradation under Thermal Irradiated Fields

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**Abstract**— This paper has presented the performance of the VCSELs in thermo-irradiated has been modeled, investigated and deeply parametrically analyzed over wide range of the affecting parameters. As well as we have processed soliton transmission capacity of VCSELs based different materials (silica-doped, aluminum gallium arsenide (AlGaAs), and polymer) in thermo-irradiated field effects. Both the ambient temperature and irradiation dose as well as the spectral wavelength possess serve reduction effects on the transmission characteristics (dispersion, and bandwidth) and consequently the transmitted bit rates and products. Useful semi-empirical relations have been cast. Thermo-irradiated penalties were computed and fitted as nonlinear relationships of useful impact in the design model for the device. Analysis of laser characteristics after irradiation showed that the main effect of radiation damage is an increase in bulk recombination that increases loss within the laser cavity. The device performance degradation is proportional to the fluence. The fluence rate is also relevant for degradation in electron irradiation. Low fluence rate leads to larger degradation compared to those associated with high fluence rate resulting from heat impact in bulk. The radiation damage of neutron is larger than irradiation damage of electron, which is caused by the difference in mass and the possibility of nuclear collision for the formation of lattice defects.

**Keywords**— VCSELs, Radiation Hardness Assurance, Reliability, Soliton Transmission and Neutrons Irradiation

## I. INTRODUCTION

Many advances have been made in laser diodes that have extended the wavelength range over which semiconductor lasers can be used, and also reduced the threshold current required for operation. Key technical advances include the use of thin heterostructures that allow efficient carrier injection over dimensions of 0.1 to 0.2  $\mu\text{m}$  [1], special means of confining carriers within the laser cavity (including interspersed oxide or semiconductor layers), strained layers that reduce losses due to Auger recombination, and small, precise lateral dimensions that provide quantum wells to limit the number of allowable states within the cavity [2].

The oldest material system is AlGaAs-GaAs, used for wavelengths between 0.65 and 0.85  $\mu\text{m}$ . Alloys of AlGaAs are lattice matched with GaAs to about 0.1 %, allowing a wide range of compositions to be used. The second material system is InGaAsP-InP. This combination also allows good lattice matching for various compositional variations [3], and is widely used for wavelengths between 1.1 and 1.6  $\mu\text{m}$  [4]. Close lattice matching is not possible for the wavelength “gap” between 0.85  $\mu\text{m}$  and 1.1  $\mu\text{m}$ . Lasers within that

region are made with InGaAs-GaAs, where there is a large amount of lattice mismatch (up to 3 %). Thus, lasers fabricated with this material contain strained layers that must be restricted in thickness in order to avoid lattice defects that would degrade efficiency and reliability. Strained layers can be made with dimensions up to approximately 100  $\mu\text{m}$  [5].

In the present study, we have investigated deeply the performance characteristics of high speed laser diodes such as vertical cavity surface emitting lasers (VCSELs) under thermal irradiated operating conditions over wide range of the affecting parameters. As well as a considerable amount of hard radiation effects studies on individual optoelectronic devices exposed to a variety of radiation conditions is deeply investigated, only little information is available on the radiation tolerance at high total dose and under neutron radiation.

## II. SCHEMATIC DIAGRAM OF AlGaAs LASER

As shown in Fig. 1, a simplified diagram of an advanced laser diode (the top and bottom electrical contacts are not shown). The structure includes double-heterojunctions that allow carrier injection over very short distances, and lateral confinement using an oxide layer to restrict carrier injection from the top contact. For this material, the difference in refractive index between the GaAs and AlGaAs regions is sufficient to confine photons to the active region, but some lasers add additional layers that surround the active region for better optical confinement. Several methods can be used for lateral confinement, including ridge structures [6]. Proton displacement damage has been studied in only a few types of laser diodes [7] and little information is available on damage in more advanced devices [8]. The

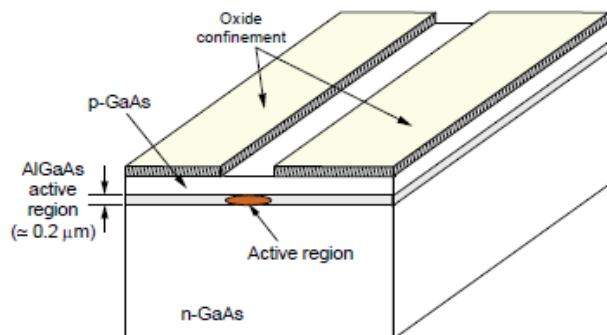


Fig. 1. Diagram of an oxide-confined AlGaAs laser

lasers represent a broad range of current laser technologies. One of the devices is a vertical cavity surface-emitting laser that provides light output from the top surface, and uses a distributed Bragg reflector in combination with lateral oxide confinement in order to obtain high efficiency and low threshold current [9].

**III. BASIC MODEL AND ANALYSIS**

The degradation of junction diode performance is known to be caused by a decrease in both the minority carrier lifetime and generation lifetime. The decrease of these lifetimes can be related to the total dose  $\phi$  by the definition of a damage coefficient k:

$$\frac{1}{t_p} = \frac{1}{t_{p0}} + k_p \phi \quad (1)$$

$$\frac{1}{t_g} = \frac{1}{t_{g0}} + k_g \phi \quad (2)$$

Where  $\phi$ ,  $k_p$ ,  $k_g$ ,  $t_{p0}$  and  $t_{g0}$  are irradiation fluence in  $n/cm^2$ , the damage coefficient of minority carrier lifetime, damage coefficient of generation lifetime, initial minority carrier lifetime and initial generation carrier lifetime, respectively. The  $t_p$  and  $t_g$  depend on the total dose and other parameters are independent of the total dose. When high energy radiation falls on a semiconductor device, energy is deposited in the semiconductor via two mechanisms; atomic displacement or ionization [10]. The relative importance of these mechanisms in a semiconductor depends on the type of radiation and the nature of the device. The electrical resistance of semiconducting metals is very much more sensitive to radiation than that of typical metals because the radiation-induced defects may markedly alter the number of electrons available for conduction. Point defects (initially produced) can be produced by electron or gamma irradiation. It can be expressed as a single displaced atom and its associated vacancy called Frankel Defect. The interaction is simply described by the number of defects/cm<sup>3</sup> created, which is given by [11]:

$$N_t = \phi a_d N_0 \quad (3)$$

Where  $a_d$  is the displacement cross-section in  $cm^2$ , and  $N_0$  is the number of lattice atoms/cm<sup>3</sup>. The point defects result in the introduction of allowed energy states within the forbidden gap of the semiconductor. This energy states lead to the following: (a) Carrier removal is the majority carrier density is reduced by the radiation fluence, (b) Mobility degradation is the mobility was found to decrease with increasing the radiation fluence, (c) Conductivity modulation since the carrier concentration and mobility both decrease with radiation, then the conductivity will also decrease, and (d) Minority carriers lifetime can be defined as the degradation rate in minority carrier lifetime and can be expressed as:

$$\frac{d\tau}{d\phi} = K_\tau \quad (4)$$

Where  $K_\tau$  is the carrier lifetime damage constant. The minority carrier lifetime  $\tau$  will be reduced by [12]:

$$\frac{1}{\tau} = \frac{1}{\tau_0} + \sum_i v_{th} \sigma_i n_i \quad (5)$$

Where  $\tau_0$  is the pre-irradiation minority carrier lifetime,  $v_{th}$  is the average thermal velocity of the minority carriers,  $n_i$  and  $\sigma_i$  are the density and capture cross section, respectively, of the recombination centers of a given type  $i$ , and the sum extends over the different types of radiation induced defects.

As long as there is no significant overlap of the defect regions produced by the individual incident particles [13], the initial defect density will be proportional to the particle fluence  $\phi$  as the following expression:

$$n_i = c_i \phi I \quad (6)$$

Where  $I$  is the light current output, and the coefficient  $c_i$  stands for the density of particular defects per unit fluence. Defining a damage constant  $K_\tau$  is:

$$K_\tau = v_{th} \sum_i c_i \sigma_i \quad (7)$$

Where  $K_\tau$  is sensitive to many factors such as material, type of radiation, particle energy, flux, device temperature and electrical bias conditions. The ratio  $\tau_0/\tau$  of the minority carrier lifetime before and after the irradiation can be related to the relative light output after irradiation assuming that the total current density in the junction is dominated by diffusion currents as the following expression [15]:

$$\left(\frac{I_0}{I}\right)^{2/3} = \frac{\tau_0}{\tau} \quad (8)$$

Where  $I_0$  is pre-irradiation light current output. Thus, the parameter  $k = \tau_0 K_\tau$  can be calculated from the relative light output (RLO) of the devices after the irradiation with a total particle fluence  $\phi$ ;

$$\left(\frac{1}{RLO}\right)^{2/3} = \left(\frac{I_0}{I}\right)^{2/3} = 1 + k \phi \quad (9)$$

The power forward current curves of different types of VCSEL devices were given in [16], with remarkable nonlinearly while in [17] it depicted in linear fashion. Based on the data of [18], the following nonlinear thermal relations for the set of the selected device were carried out:

$$P(I,T) = p_0 + p_1 I + p_2 I^2, \quad \text{mWatt} \quad (10)$$

Where the set of parameters  $\{p_0, p_1, \text{ and } p_2\}$  are polynomial functions of ambient temperature  $T$ .

$$p_0 = 0.73 - 0.00169 T + 0.000345 T^2, \quad (11)$$

$$p_1 = 2.5 - 0.0072 T + 0.0002 T^2, \quad (12)$$

$$p_2 = -7.3 - 0.002 T + 0.000065 T^2, \quad (13)$$

The power of VCSEL under the effects of irradiation can be expressed in the following formula:

$$P(I,T,\phi) = P(I,T) F_p(\phi) \quad (14)$$

Where  $F_p(\phi)$  is a function of the irradiation fluence  $\phi$ , can be expressed as follows:

$$F_p(\phi) = 1 + \alpha_1 \phi + \alpha_2 \phi^2 \quad (15)$$

Where the set of the coefficients of  $\alpha_1=0.0005$ ,  $\alpha_2=-0.001$ . The standard single mode fiber link within VCSEL based silica doped material which the investigation of the spectral variations of the waveguide refractive-index,  $n$  require empirical equation under the form [19]:

$$n^2 = 1 + \frac{A_1 \lambda^2}{\lambda^2 - A_2^2} + \frac{A_3 \lambda^2}{\lambda^2 - A_4^2} + \frac{A_5 \lambda^2}{\lambda^2 - A_6^2} \quad (16)$$

The empirical equation coefficients as a function of temperature and Germania mole fraction,  $x$  as:

$A_{1S}=0.691663+0.1107001x$ ,  
 $A_{2S}=(0.0684043+0.000568306x)^2(T/T_0)^2$ ,  
 $A_{3S}=0.4079426+0.31021588x$ ,  
 $A_{4S}=(0.1162414+0.03772465x)^2(T/T_0)^2$ ,  $A_{5S}=0.8974749-0.043311091x$ ,  $A_{6S}=(9.896161+1.94577x)^2$ . Where  $T$  is ambient temperature in K, and  $T_0$  is the room temperature and is considered as 300 K. First and second differentiation of empirical equation w. r. t operating wavelength  $\lambda$  as in

Ref. [20] for both silica-doped and polymer materials. For the plastic fiber material, the coefficients of the Sellmeier equation and refractive-index variation with ambient temperature are given as:  $A_{1p}=0.4963$ ,  $A_{2p}=0.6965$  ( $T/T_0$ ),  $A_{3p}=0.3223$ ,  $A_{4p}=0.718$  ( $T/T_0$ ),  $A_{5p}=0.1174$ , and  $A_{6p}=9.237$ . For Aluminum Gallium Arsenide (AlGaAs), the set of parameters required to characterize the temperature and operating wavelength dependence of the refractive-index, where empirical equation is under the form of [20]:

$$n^2 = B_1 + \frac{B_2}{\lambda^2 - B_3} - B_4\lambda^2, \quad (17)$$

The set of parameters is recast and dimensionally adjusted as:  $B_1=10.906$ ,  $B_2=0.975$ ,  $B_3=0.3464499$  ( $T/T_0$ ),  $B_4=0.002467$  ( $0.93721+2.0857 \times 10^{-4} T$ ). Where  $T$  is the ambient temperature in K, and  $T_0$  is the room temperature (300 K). We differentiate of previous empirical equation w. r. t  $\lambda$  as mentioned in Ref. [19].

The idea of soliton transmission is to guide the nonlinearity to the desired direction and use it for our benefit. When soliton pulses are used as an information carrier, the effects of dispersion and nonlinearity balance each other and thus don't degrade the signal quality with the propagation distance. In addition, the unique features of soliton transmission can help to solve the problems of data transmission, because the soliton data looks essentially the same at different distances along the transmission, the soliton type of transmission is especially attractive for all-optical data networking. Moreover, because of the high quality of the pulses and return-to-zero (RZ) nature of the data the soliton data is suitable for all-optical processing. In any infinitesimal segment of fiber, dispersion on one hand and non linearity of the refractive-index on the other hand produce infinitesimal modulation angles which exactly compensate reciprocally. In the sense that their sum is an irrelevant constant phase shift. Under such conditions the pulse shape is the same everywhere. The soliton waveform be used with a peak power under the effects of irradiation fluence, temperature and threshold current [21]:

$$P(I, T, \varphi) = \frac{\lambda_s^3 D_t A_{eff}}{12.7 c n_{nl} \Delta \tau^2}, \quad (18)$$

Where  $n_{nl}$  is the nonlinear Kerr coefficient,  $2.6 \times 10^{-20}$   $m^2/Watt$ ,  $\lambda$  is the operating signal wavelength in  $\mu m$ ,  $A_{eff}$  is the effective area of the plastic fiber in  $\mu m^2$ ,  $c$  is the velocity of light ( $3 \times 10^8$  m/sec). The total bandwidth is based on the total chromatic dispersion for silica-doped and AlGaAs based VCSEL is given by:

$$D_t = D_m + D_w \quad (19)$$

Where  $D_m$  is the material dispersion and  $D_w$  is the waveguide dispersion as mentioned in Ref. [22]. Also the total dispersion coefficient in  $nsec/m^2$  of the polymer material based VCSEL is given by:

$$D_t = (M_{md} + P), \quad (20)$$

In which  $M_{md}$  is the material dispersion and  $P$  is the profile dispersion as listed in Ref. [22]. Then the total pulse broadening from Eq. 19 can be expressed as the following:

$$\Delta \tau = \sqrt{\frac{\lambda^3 D_t A_{eff}}{12.7 P(I, T, \varphi) n_{nl} c}}, \quad (21)$$

The bandwidth for standard single mode fibers for both materials based VCSEL device is given by [23]:

$$BW_{sig} = \frac{0.44}{\Delta \tau \cdot L}, \quad (22)$$

Then the soliton transmission bit rate is given by [23]:

$$B_{RS} = \frac{1}{10 \Delta \tau} = \frac{0.1}{\Delta \tau}, \quad (23)$$

Where  $L$  is the transmission VCSEL laser diode reach.

#### IV. SIMULATION RESULTS AND DISCUSSIONS

We have investigated transmission bit rates and bandwidth transmission reach products of VCSEL devices under the effects of the thermo-Irradiated environments under the assumed set of the operating parameters and ranges as: irradiation fluence  $2 \times 10^{14} \leq \varphi \leq 5 \times 10^{14}$   $n/cm^2$ , Ambient temperature  $290 \leq T, K \leq 330$ , applied drive current  $10 \leq I, mA \leq 50$ , room temperature  $T_0=300$  K, effective area  $A_{eff}=85 \mu m^2$ , index exponent  $\alpha=2$ , spectral linewidth of the optical source  $\Delta \lambda=0.2$  nm, signal power  $P_{so}=2$  mWatt, germanium dopant ratio  $x=0.3$ , Operating optical signal wavelength  $0.65 \leq \lambda_s, \mu m \leq 1.65$ , silica relative refractive index difference  $\Delta n_s=0.003$ , AlGaAs relative refractive index difference  $\Delta n_{AlGaAs}=0.003$ , polymer relative refractive index difference  $\Delta n_p=0.03$ , initial device gain,  $g_0=8.5 \times 10^6$ , thermal rate  $R_{th}=0.85$ , transmission laser diode reach  $L=50$  m,  $N_0=1.3 \times 10^6$ . It is necessary to stress that the major generator of features is the irradiation fluence due to the damage which causes dislocations; and radiation-induced defects is the band gap. Based on the above model and the series of the operating parameters variations of the set of five causes {3-dB signal transmitted bandwidth  $f_{3-dB}$ , device transmission bit rates and products} against the variations of a set of effects {irradiation fluence  $\varphi$ , ambient temperature  $T$ , applied threshold current  $I$ } are displayed in the series of Figs. (2-14):

- i) As shown in Fig. 2 has assured that as the drive current of the VCSEL diode increases, this leads to increase of output device power. We have observed that after two neutrons fluences, the output device power is decreased compared to that its value before irradiation.
- ii) As shown in the series of Figs. (3-8), as ambient temperature increases, these results in decreasing in soliton transmission bit rates for different materials based VCSEL devices after or before irradiation. As well as we have demonstrated after two neutrons fluences, the soliton transmission bit rates for different based VCSEL devices degraded and decreased fastly. This illustrates the harmful effects of the neutrons irradiation on the performance characteristics of these devices, especially for transmission bit rates.
- iii) As shown in the series of Figs. (9-14), as ambient temperature increases, these results in decreasing in transmitted signal bandwidth for different materials based VCSEL devices after or before irradiation. As well as we have indicated after two neutrons fluences, the transmitted signal bandwidth for different based VCSEL devices degraded and decreased fastly. This illustrates the harmful effects of the neutrons irradiation on the performance characteristics of these devices, especially for transmitted signal bandwidths.
- iv) As shown in the series of Figs. (2-14), we have assured that as the drive current increases, this leads to increase of output device power, and then to increase both transmission bit rates and transmitted signal bandwidths.

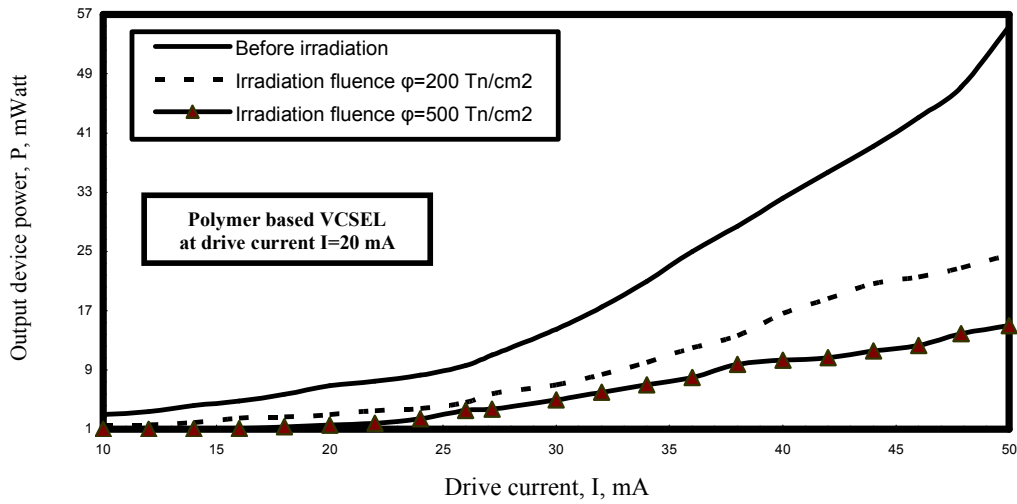


Fig. 2. Variations of the output power against drive current for a typical VCSEL after two neutrons fluences at the assumed set of parameters.

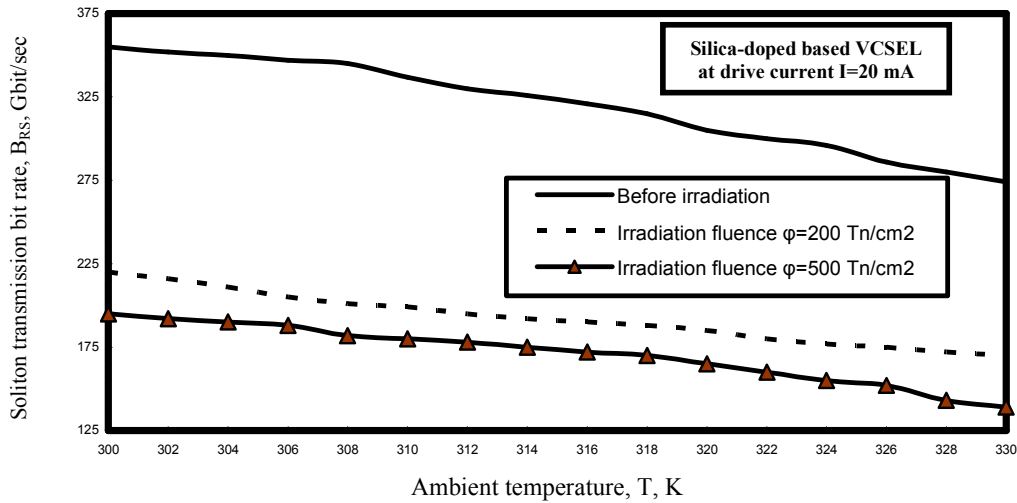


Fig. 3. Variations of the soliton transmission bit rate against ambient temperature after two neutrons fluences at the assumed set of parameters.

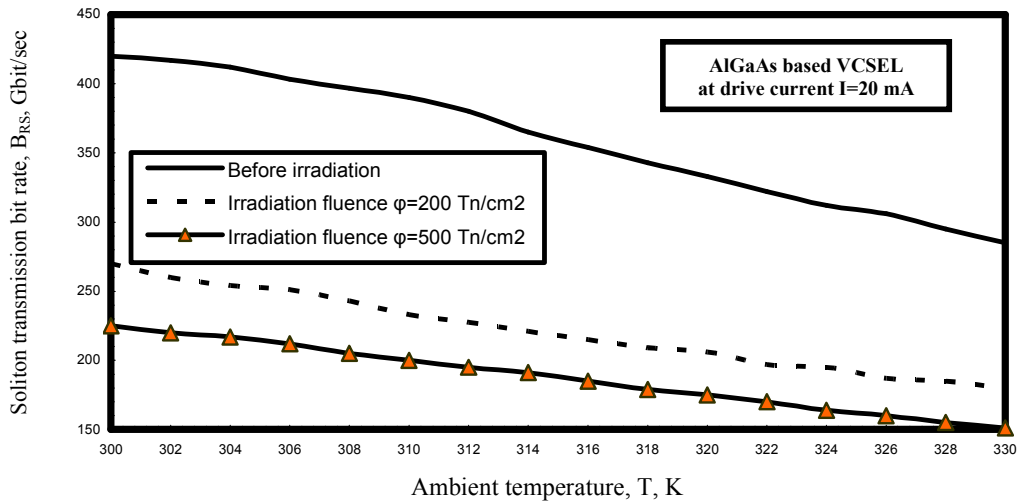


Fig. 4. Variations of the soliton transmission bit rate against ambient temperature after two neutrons fluences at the assumed set of parameters.

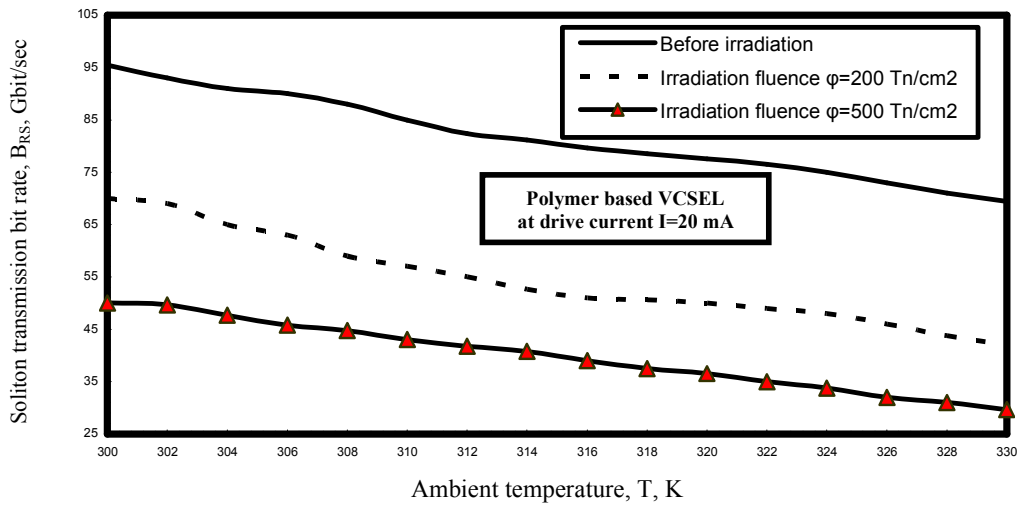


Fig. 5. Variations of the soliton transmission bit rate against ambient temperature after two neutrons fluences at the assumed set of parameters.

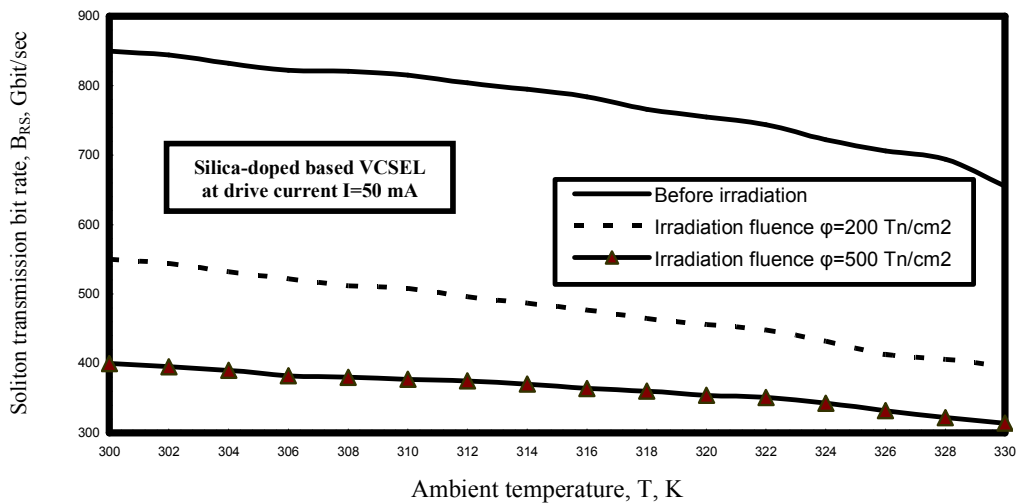


Fig. 6. Variations of the soliton transmission bit rate against ambient temperature after two neutrons fluences at the assumed set of parameters.

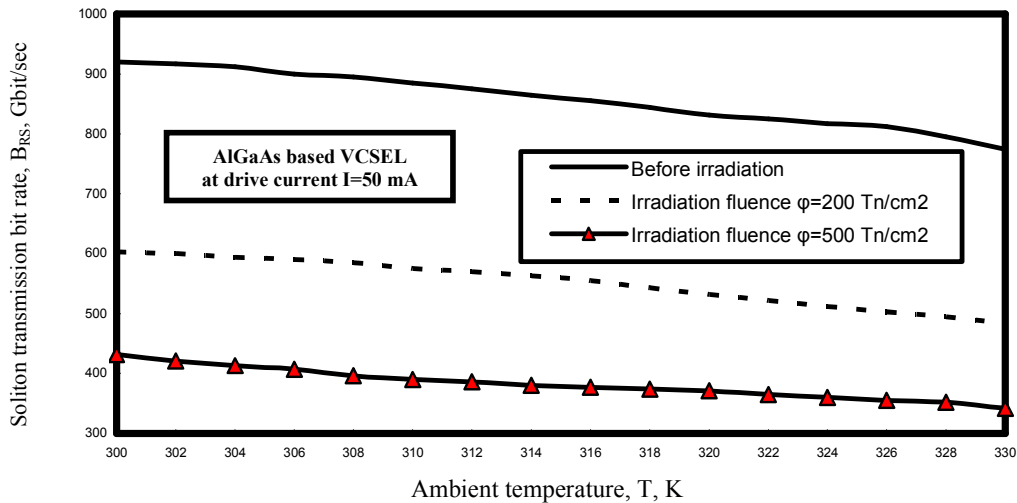


Fig. 7. Variations of the soliton transmission bit rate against ambient temperature after two neutrons fluences at the assumed set of parameters.

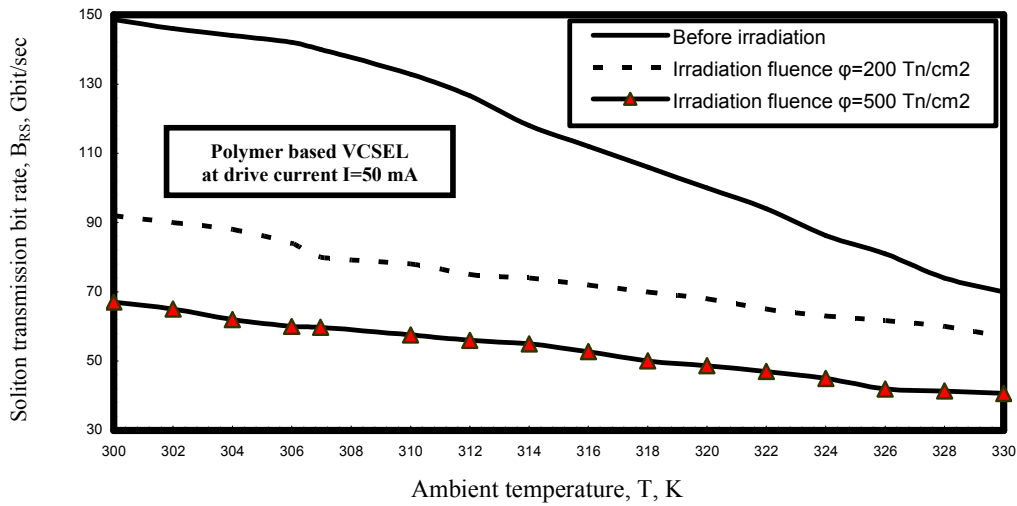


Fig. 8. Variations of the soliton transmission bit rate against ambient temperature after two neutrons fluences at the assumed set of parameters.

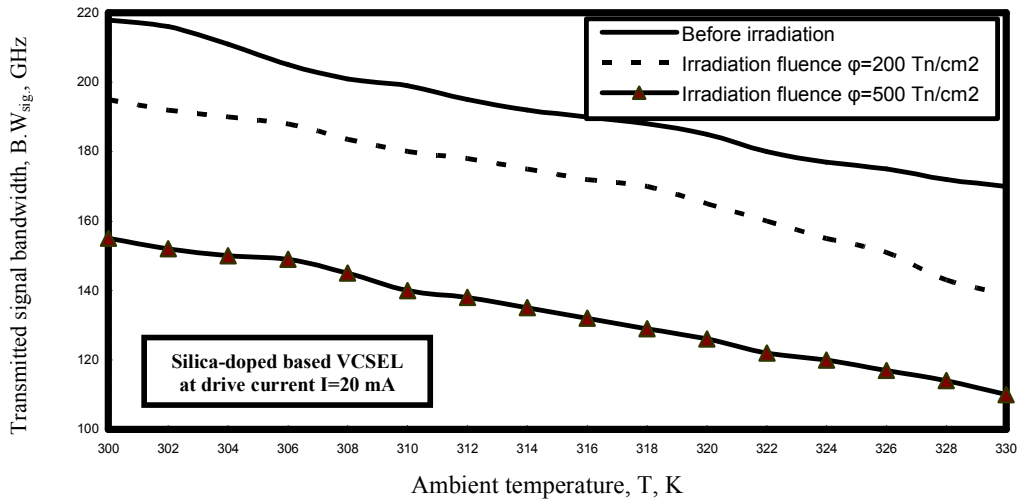


Fig. 9. Variations of the transmitted signal bandwidth against ambient temperature after two neutrons fluences at the assumed set of parameters.

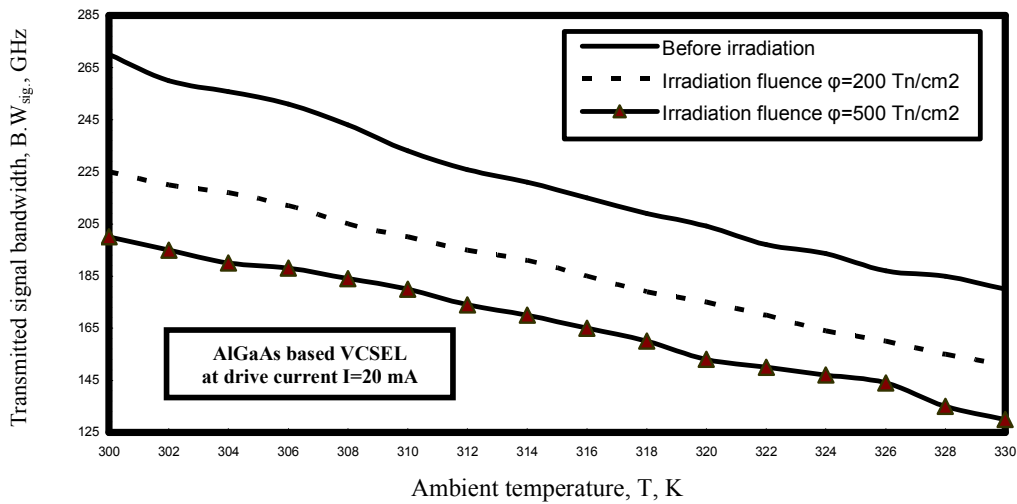


Fig. 10. Variations of the transmitted signal bandwidth against ambient temperature after two neutrons fluences at the assumed set of parameters.

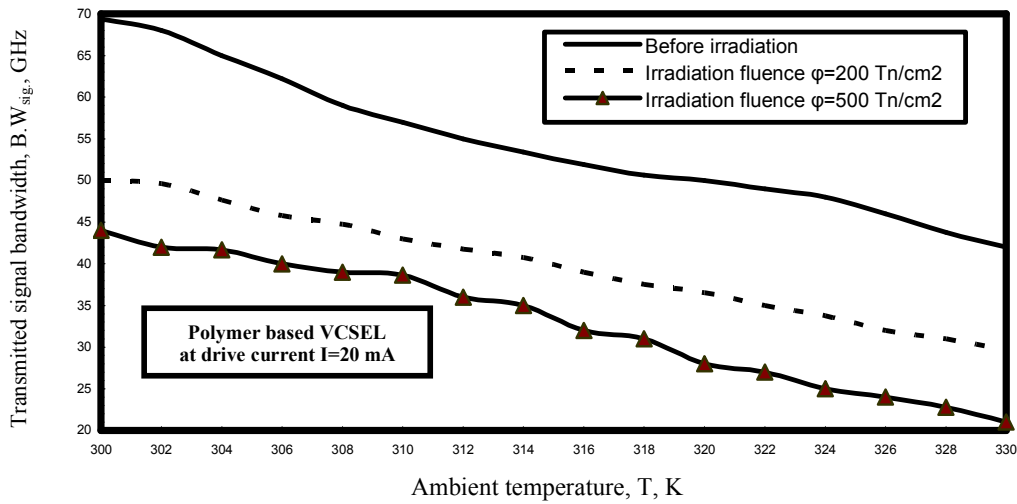


Fig. 11. Variations of the transmitted signal bandwidth against ambient temperature after two neutrons fluences at the assumed set of parameters.

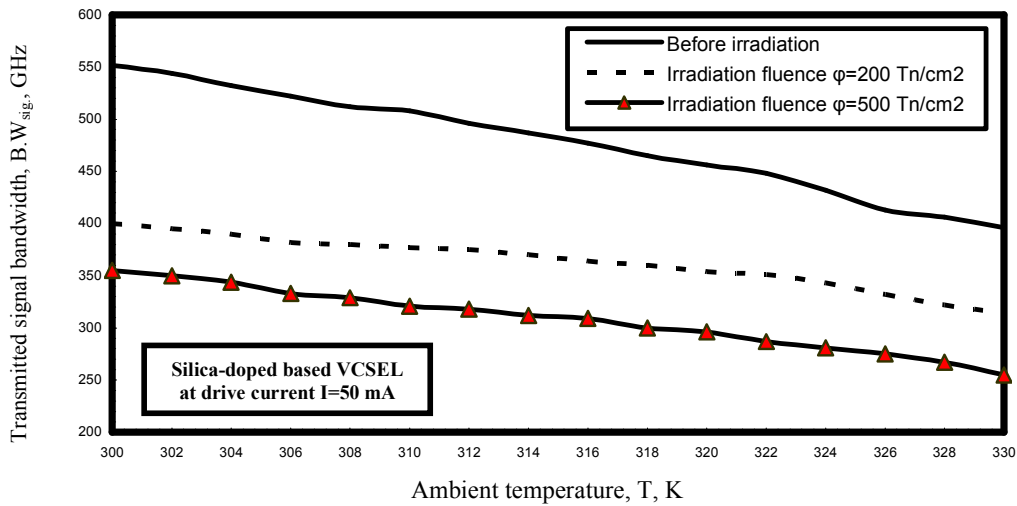


Fig. 12. Variations of the transmitted signal bandwidth against ambient temperature after two neutrons fluences at the assumed set of parameters.

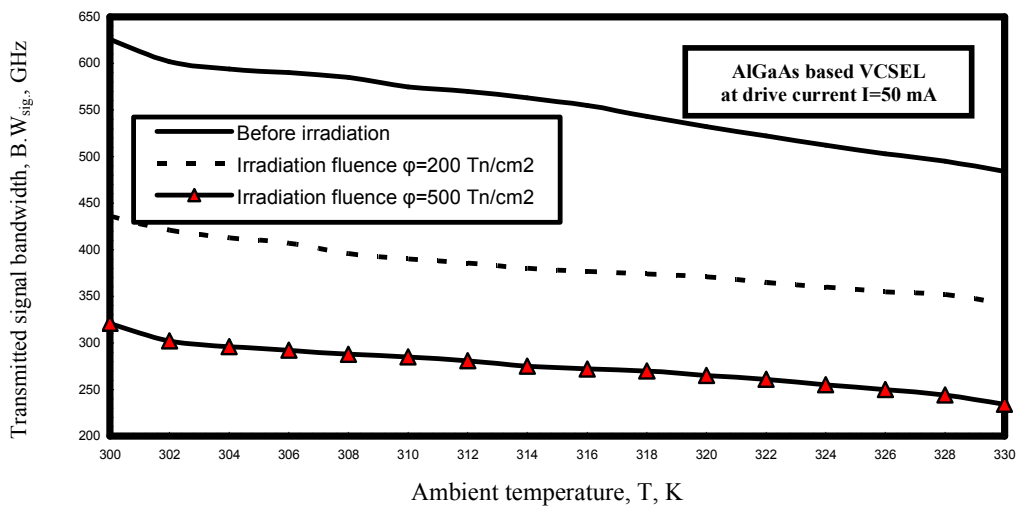


Fig. 13. Variations of the transmitted signal bandwidth against ambient temperature after two neutrons fluences at the assumed set of parameters.

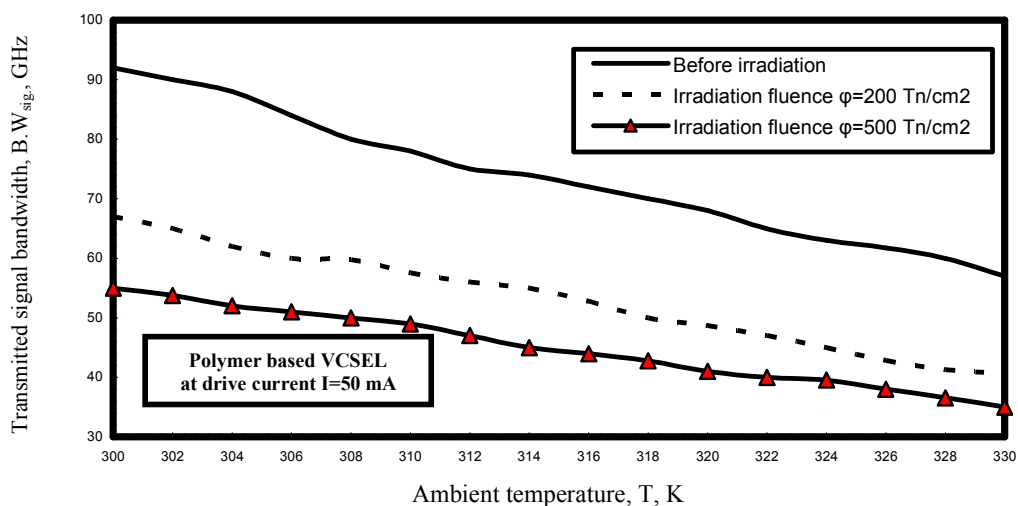


Fig. 14. Variations of the transmitted signal bandwidth against ambient temperature after two neutrons fluences at the assumed set of parameters.

## V. CONCLUSIONS

In a summary, we have deeply investigated the low performance characteristics of VCSEL devices under thermal irradiation fields over wide range of the affecting parameters. It is theoretically found that the increased drive current of VCSEL devices this results in the increased output device power. As well as we have observed that the output device power is decreased dramatically after two neutrons fluences. It is also indicated that the increased both ambient temperature and neutrons irradiation fluences this leads to the decreased of both soliton transmission bit rates and transmitted signal bandwidths for different materials based VCSEL devices. Therefore we have demonstrated the harmful effects of both ambient temperature and neutrons irradiation fluences over different materials based VCSEL devices that affect negatively on the performance characteristics of the VCSEL devices.

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